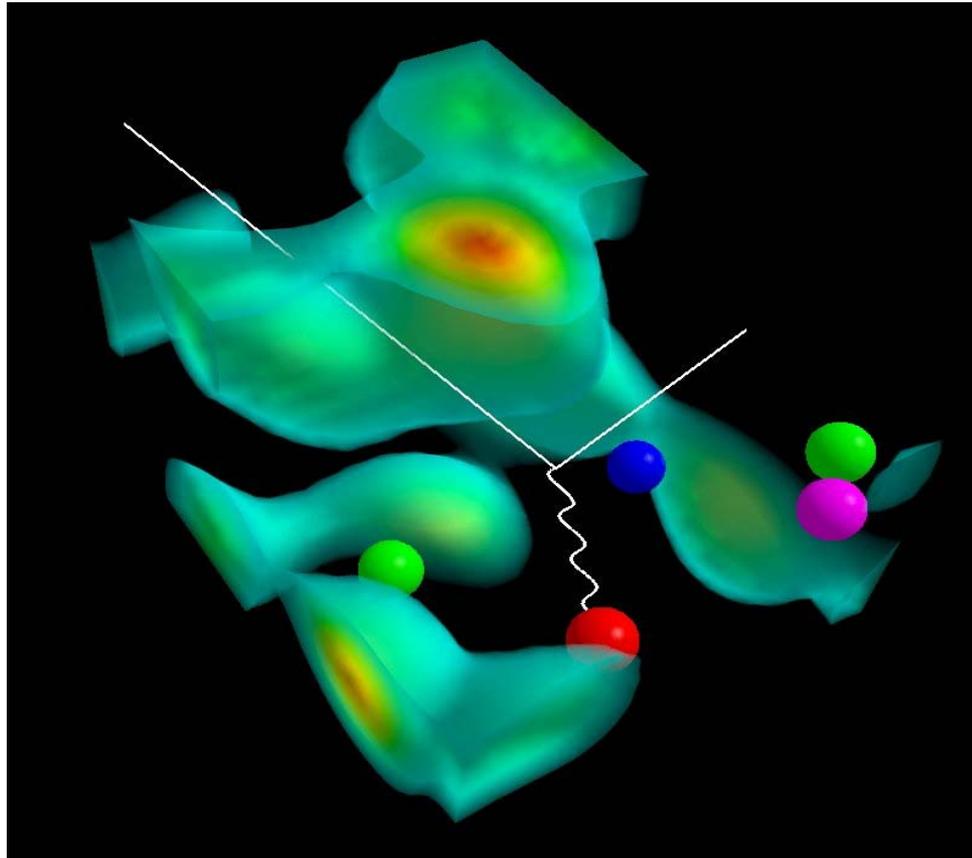


Strangeness in the Nucleon - Introduction



Anthony W. Thomas

CIPANP San Diego : May 30th 2009



Operated by Jefferson Science Associates for the U.S. Department of Energy

Thomas Jefferson National Accelerator Facility



Outline

- **Strangeness contribution is a vacuum polarization effect, analogous to Lamb shift in QED**

Hydrogen Atom, Electron (g-2)-factor, QED

$$g_e = 2 \left(1 + \frac{\alpha}{2\pi} - 0.328 \frac{\alpha^2}{\pi^2} + \dots \right)$$

- **It is a fundamental test of non-perturbative QCD**

Outline (cont.)

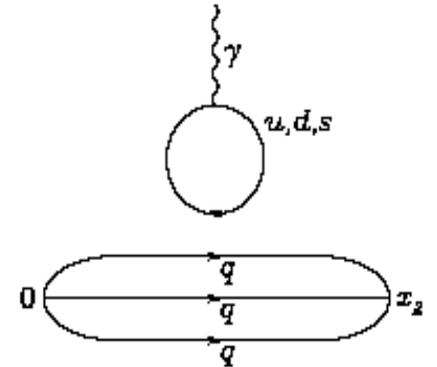
There have been a number of major steps forward recently, both theory and experiment

- Contribution to nucleon mass
- Corresponding sigma term (Young, yesterday)
- Calculation of $G_{E,M}^s(Q^2)$:
 - Direct: Kentucky (χ QCD : K.-F. Liu)
 - Indirect: JLab-Adelaide
- Experimental determination of $G_{E,M}^s(Q^2)$
 - G0 (Beise, this session); Mainz PVA4 ([arXiv:0903.2733](https://arxiv.org/abs/0903.2733)); Happex and Bates
- Agreement between theory and experiment excellent
 - consistent global analysis valuable

Strangeness & Electromagnetic Form Factors

Experiment: Need Parity Violation

Theory: “Disconnected diagram”

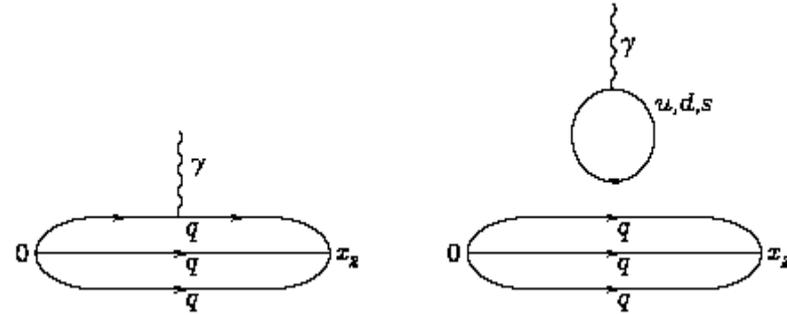


Hydrogen Atom, Electron (g-2)-factor, QED

$$g_e = 2 \left(1 + \frac{\alpha}{2\pi} - 0.328 \frac{\alpha^2}{\pi^2} + \dots \right)$$

Magnetic Moments within QCD

Leinweber and Thomas, Phys Rev D62 (2000)



CS $\left\{ \begin{array}{l} \mathbf{p} = 2/3 \mathbf{u}^p - 1/3 \mathbf{d}^p + \mathbf{O}_N \\ \mathbf{n} = -1/3 \mathbf{u}^p + 2/3 \mathbf{d}^p + \mathbf{O}_N \end{array} \right.$



$$2\mathbf{p} + \mathbf{n} = \mathbf{u}^p + 3 \mathbf{O}_N$$

(and $\mathbf{p} + 2\mathbf{n} = \mathbf{d}^p + 3 \mathbf{O}_N$)

$\left\{ \begin{array}{l} \Sigma^+ = 2/3 \mathbf{u}^\Sigma - 1/3 \mathbf{s}^\Sigma + \mathbf{O}_\Sigma \\ \Sigma^- = -1/3 \mathbf{u}^\Sigma - 1/3 \mathbf{s}^\Sigma + \mathbf{O}_\Sigma \end{array} \right.$



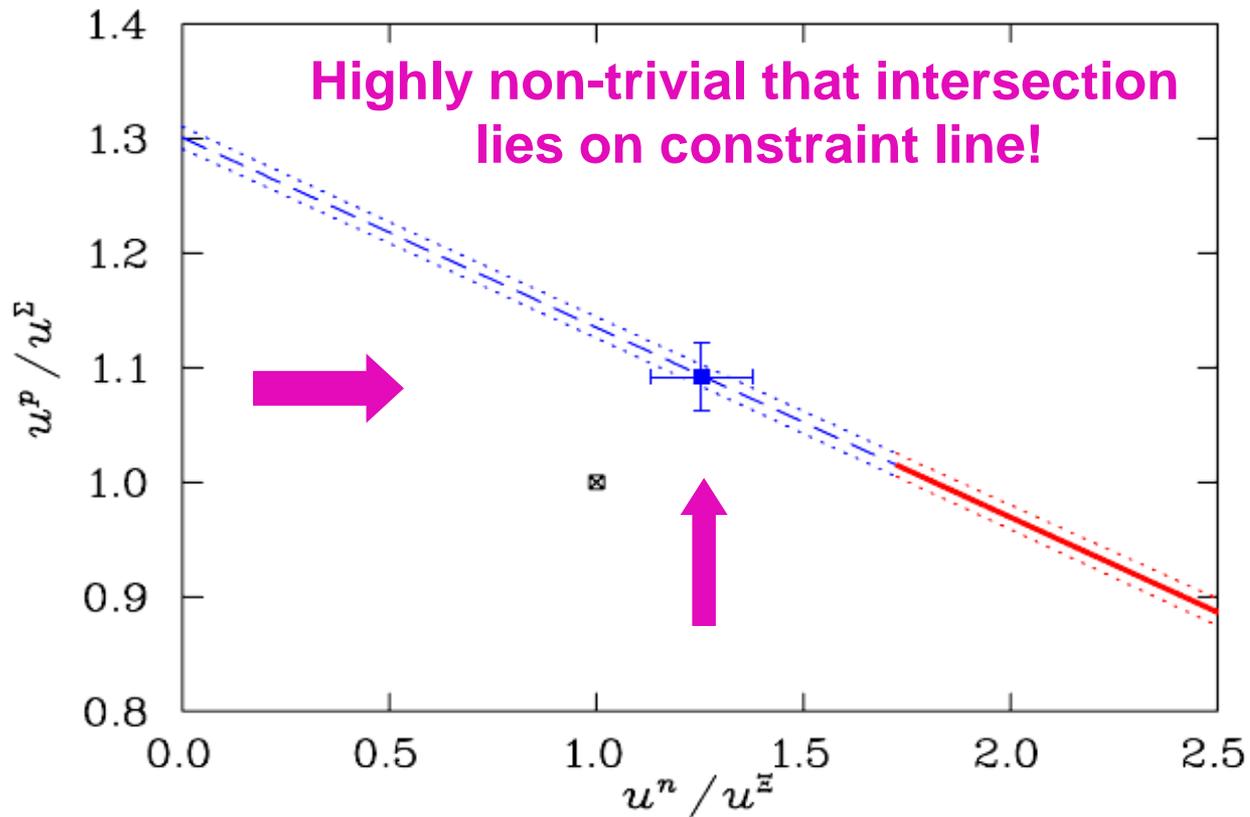
$$\Sigma^+ - \Sigma^- = \mathbf{u}^\Sigma$$

HENCE: $\mathbf{O}_N = 1/3 [2\mathbf{p} + \mathbf{n} - (\mathbf{u}^p / \mathbf{u}^\Sigma) (\Sigma^+ - \Sigma^-)]$

Just these ratios from Lattice QCD

$$\mathbf{O}_N = 1/3 [\mathbf{n} + 2\mathbf{p} - (\mathbf{u}^n / \mathbf{u}^\Xi) (\Xi^0 - \Xi^-)]$$

Accurate Final Result for G_M^s



$1.25 \quad 0.12$

Yields : $G_M^s = -0.046 \quad 0.019 \mu_N$

Leinweber et al., (PRL June '05) hep-lat/0406002

State of the Art Magnetic Moments

	QQCD	Valence	Full QCD	Expt.
p	2.69 (16)	2.94 (15)	2.86 (15)	2.79
n	-1.72 (10)	-1.83 (10)	-1.91 (10)	-1.91
Σ^+	2.37 (11)	2.61 (10)	2.52 (10)	2.46 (10)
Σ^-	-0.95 (05)	-1.08 (05)	-1.17 (05)	-1.16 (03)
Λ	-0.57 (03)	-0.61 (03)	-0.63 (03)	-0.613 (4)
Ξ^0	-1.16 (04)	-1.26 (04)	-1.28 (04)	-1.25 (01)
Ξ^-	-0.65 (02)	-0.68 (02)	-0.70 (02)	-0.651 (03)
u^p	1.66 (08)	1.85 (07)	1.85 (07)	1.81 (06)
u^Ξ	-0.51 (04)	-0.58 (04)	-0.58 (04)	-0.60 (01)

January 2006: G_E^s by same technique

In this case only know Σ^- radius (and p and n)

$$2p + n = u^p + 3 O_N \qquad p + 2n = d^p + 3 O_N$$

$$) \quad \langle r^2 \rangle_s = 0.000 \pm 0.006 \pm 0.007 \text{ fm}^2 ; 0.002 \pm 0.004 \pm 0.004 \text{ fm}^2$$

(c.f. using Σ^- : $-0.007 \pm 0.004 \pm 0.007 \pm 0.021 \text{ fm}^2$)

$$G_E^s(0.1 \text{ GeV}^2) = +0.001 \pm 0.004 \pm 0.004$$

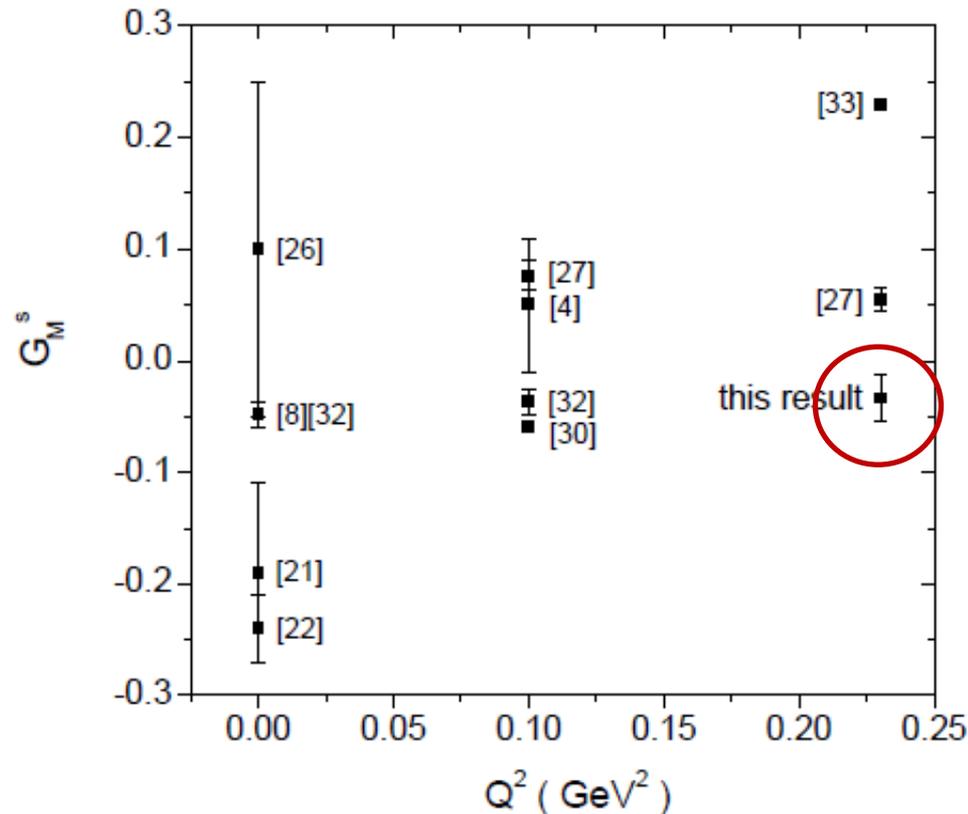
(up to order Q^4)

Note consistency and level of precision!

Leinweber, Young et al., hep-lat/0601025: Jan 2006



Indirect lattice calculation at $Q^2 = 0.23 \text{ GeV}^2$

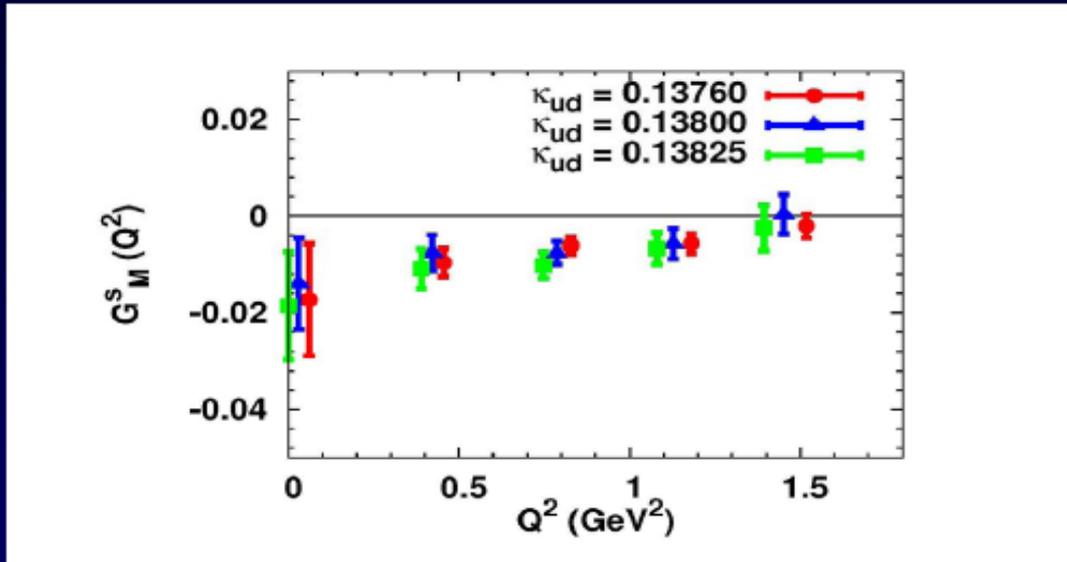


$-0.034 \pm 0.031 \mu_N$

Wang et al. : arXiv:0807.0944 (PR C in press)

Direct Calculation of $G_M^S(Q^2)$ – K.-F. Liu et al.

Strangeness Magnetic Form Factors with 3 Quark Masses
 ($m_n = 0.6, 0.7, 0.8$ GeV); T. Doi et al. (χ QCD) arXiv:0903.3232



$$G_M^S(Q^2 = 0) = -0.017(25)(07) \mu_N$$

c.f. -0.046 ± 0.019 (Leinweber et al.)

N.B. Expect increase of order 1.8 when light quark mass \rightarrow physical value with m_s fixed (Wang et al., hep-ph/0701082 :Phys Rev D75, 2008)

Moments of Strange Parton Distribution and Strangeness Magnetic Moment

- Hadronic Tensor in Euclidean Path-Integral Formalism
- $\langle x \rangle_s$ and $\langle x \rangle_{u+d}$ (D.I.)
- $\langle x^2 \rangle_s$
- Glue momentum fraction
- Strangeness Magnetic Moment

χ QCD Collaboration:

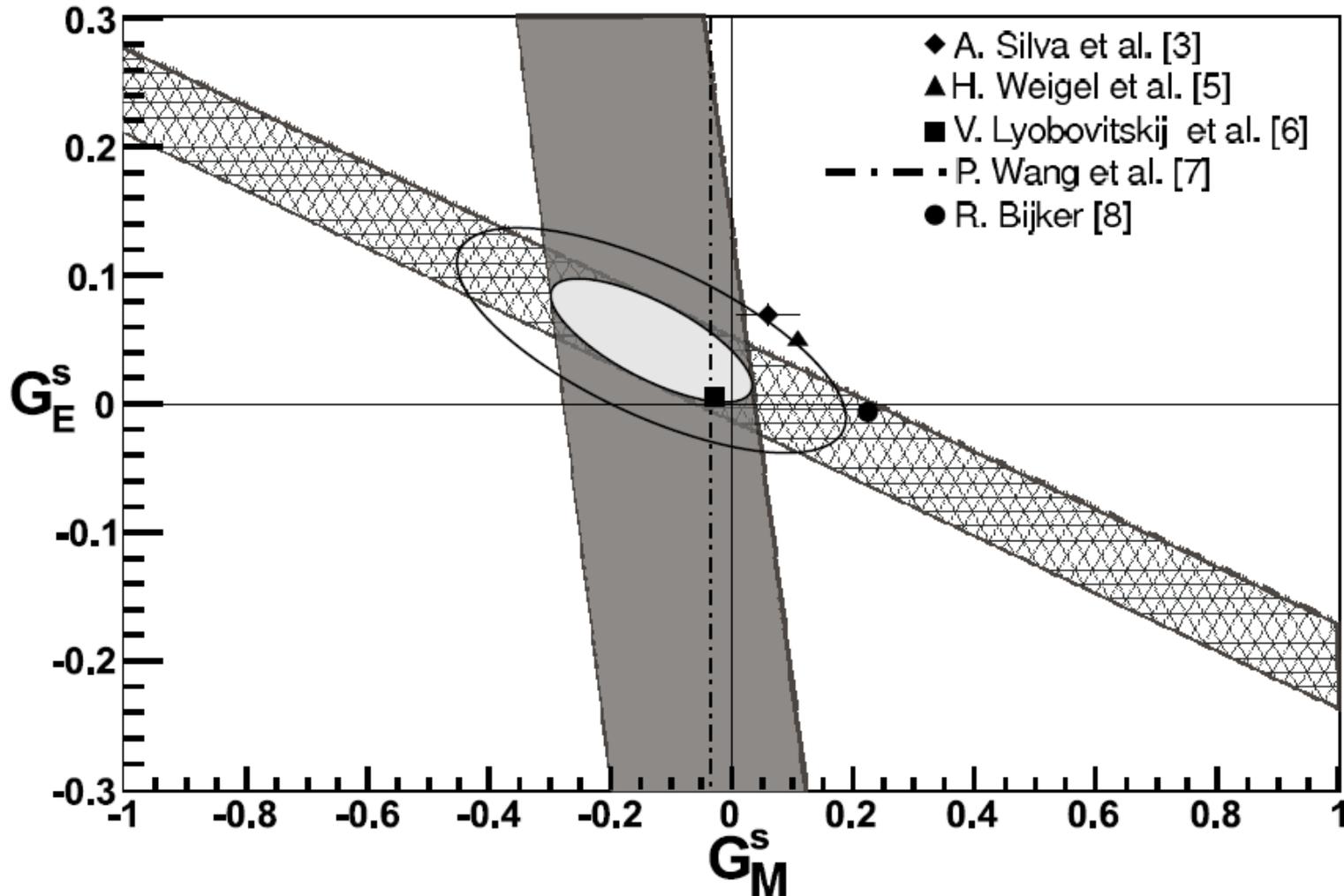
A. Alexandru, Y. Chen, T. Doi, S.J. Dong, T. Draper, I. Horvath, B. Joo, F. Lee, A. Li, K.F. Liu, N. Mathur, T. Streuer, H. Thacker, J.B. Zhang



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PVA4 2009: $Q^2 = 0.22 \text{ GeV}^2$

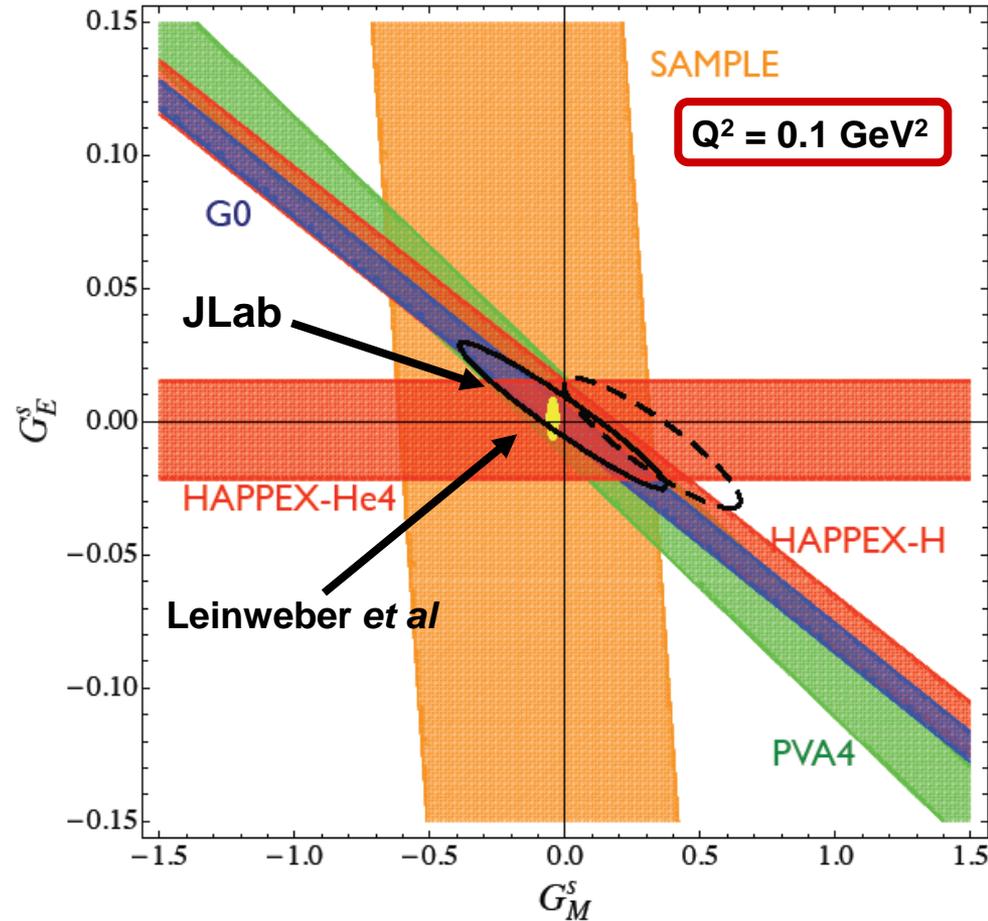
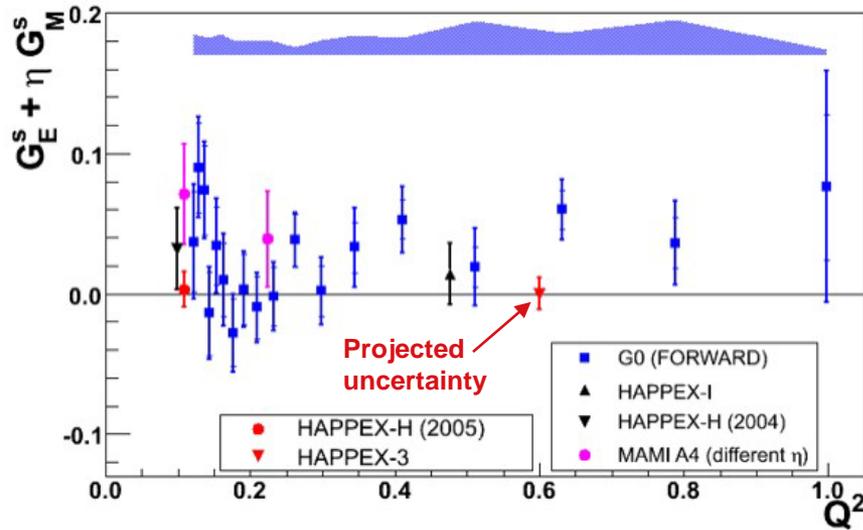
arXiv: 0903.2733v1



$$G_M^s = -0.14 \pm 0.11 \pm 0.11 \mu_N ; G_E^s = 0.050 \pm 0.038 \pm 0.019$$

Global Analysis of PVES Data

From NSAC Long Range Plan 2007:



- Proton not all that strange
- Separation possible at 0.1 GeV²
- New data coming at 0.23 and 0.6 GeV² (PVA4, G0, HAPPEX III)

Global analysis: Young et al., PRL 99 (2007)122003



“Back of the Envelope” Estimates

Nowhere that current quark masses enter dynamics
- always constituent quark masses

- Hence s-sbar pair costs 1.0-1.1 GeV plus KE
- K - Λ costs 0.65 GeV plus KE (and coupling $\gg \pi N$)
(K- Σ much smaller) ignore)
- Lots of evidence that $P_{\pi N} \sim 20\%$) $P_{K\Lambda} \sim 5\%$

$$G_M^s \sim - 3 P_{K\Lambda} [2/3 (+0.61 + 1/3) + 1/3(-0.61 + 0)]$$

$$\approx - 0.07 \mu_N$$

Thomas & Young, nucl-th/0509082



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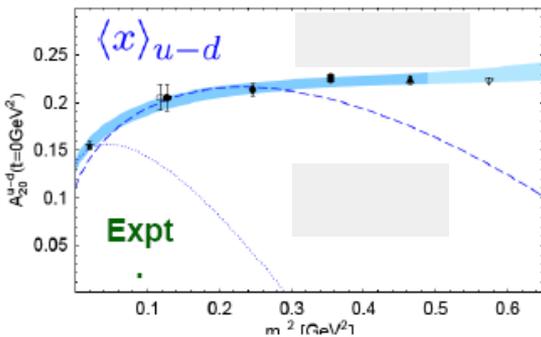
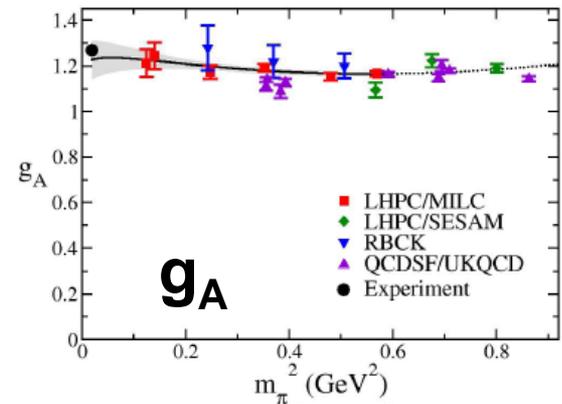
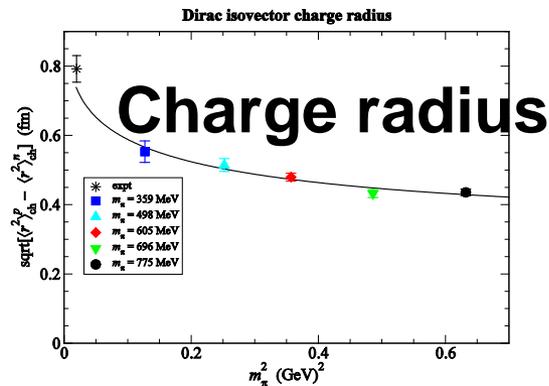
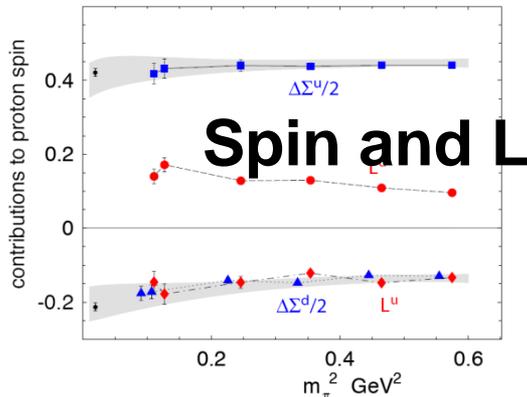
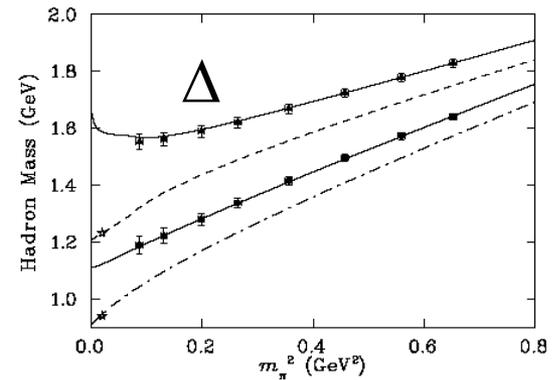
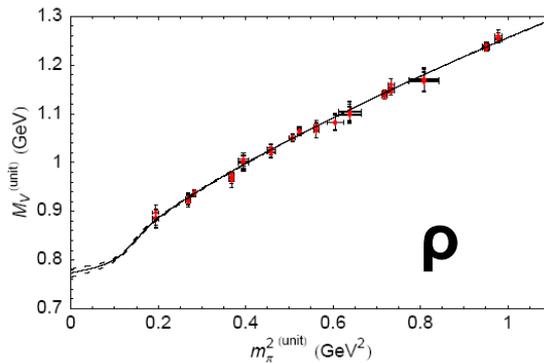
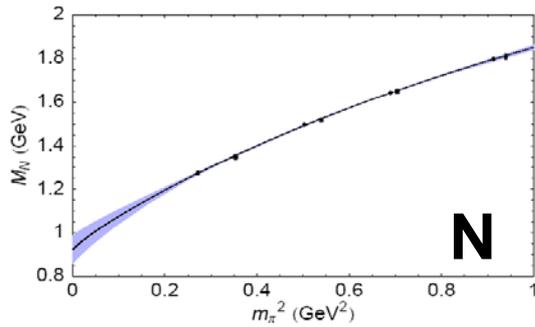
Nucleon Mass

- How much do strange quarks contribute to the mass of the nucleon ?
- Typical value is ~ 300 MeV – $1/3^{\text{rd}}$ of M_N
 - e.g. Nelson & Kaplan, 335 ± 132 MeV
(**Phys. Lett. B192, 193 (1987)**)
 - or 113 ± 108 MeV, Borasoy & Meissner (1997)
(**hep-ph/9711361**)
- Of importance for kaon condensation in dense matter
and in search for dark matter

Can now address this issue with lattice QCD data

- In fact study available data on whole octet as a function of m_π and m_K
- Data does not lie in “power counting region”
 - hence use finite range regularization (FRR)
- FRR suppresses Goldstone boson loops when Compton wavelength is too small ($m_{GB} > 0.4 \text{ GeV}$)

The "big picture"



Is it believable that smooth behavior for m_π above 400 MeV is a result of a different accidental cancellation in every case??

$$a + b m_\pi^2 + c m_\pi^3 + d m_\pi^4 \ln m_\pi + m_\pi^5 + \dots$$

Octet-baryon masses

- Leading-order expansion $O(1)$

$$M_N = M_0 + 2(\alpha_M + \beta_M)m_q + 2\sigma_M(2m_q + m_s)$$

$$M_\Lambda = M_0 + (\alpha_M + 2\beta_M)m_q + \alpha_M m_s + 2\sigma_M(2m_q + m_s)$$

$$M_\Sigma = M_0 + \frac{1}{3}(5\alpha_M + 2\beta_M)m_q + \frac{1}{3}(\alpha_M + 4\beta_M)m_s + 2\sigma_M(2m_q + m_s)$$

$$M_\Xi = M_0 + \frac{1}{3}(\alpha_M + 4\beta_M)m_q + \frac{1}{3}(5\alpha_M + 2\beta_M)m_s + 2\sigma_M(2m_q + m_s)$$

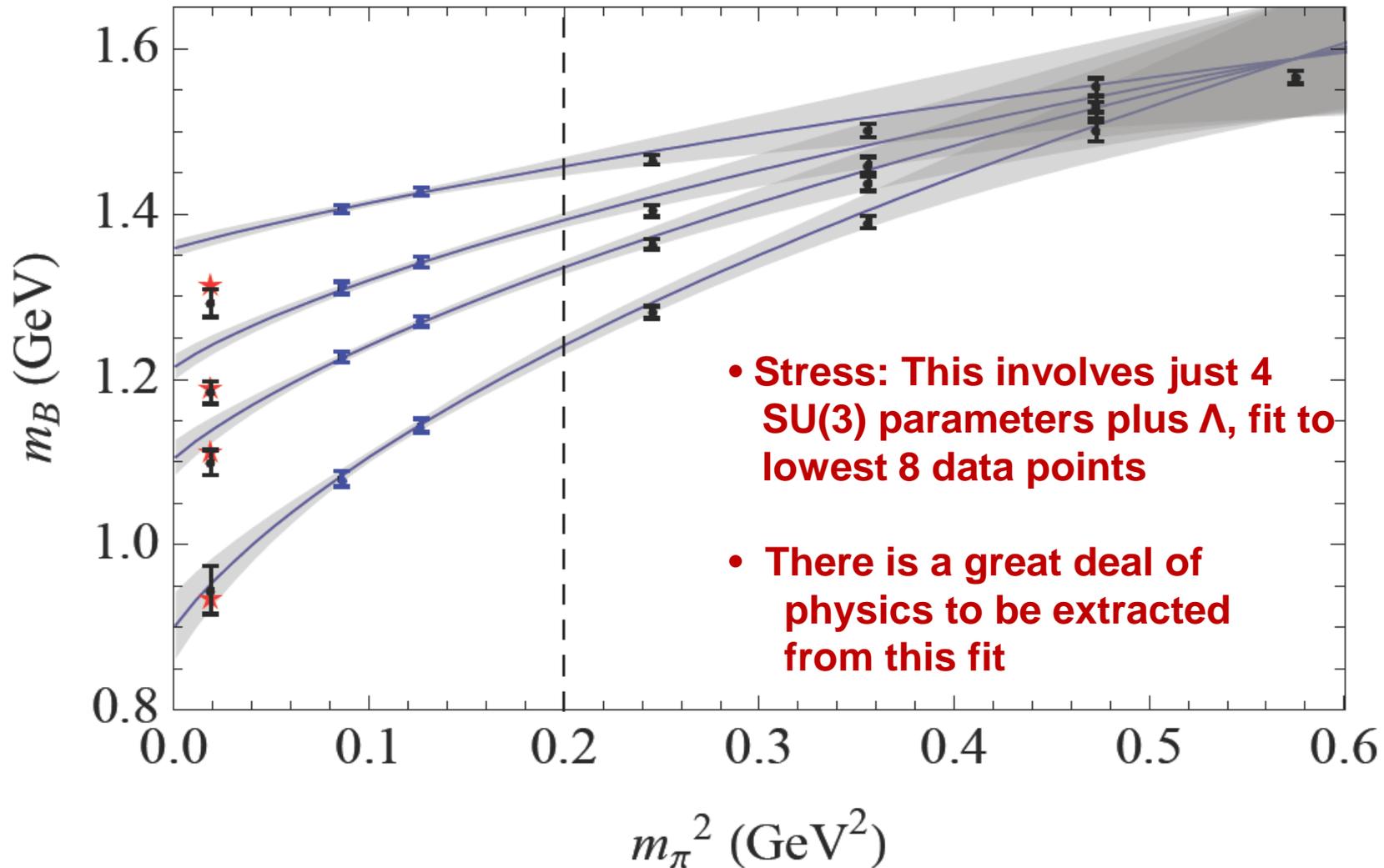
$$m_\pi^2 = 2Bm_q \quad m_K^2 = B(m_q + m_s)$$

$$m_q \rightarrow \frac{m_\pi^2}{2B}, \quad m_s \rightarrow \frac{2m_K^2 - m_\pi^2}{2B} \quad \{\alpha, \beta, \sigma\} \rightarrow B\{\alpha', \beta', \sigma'\}$$

Fit using SU(3) expansions plus FRR loops (π , η and K)

LHPC Data

(Walker-Loud et al., arXiv:0806.4549)



Young & Thomas, arXiv:0901.3559 [nucl-th]

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Summary of Results of Combined Fits

(of 2008 LHPC & PACS-CS data)

B	Mass (GeV)	Expt.	$\bar{\sigma}_{Bl}$	$\bar{\sigma}_{Bs}$
N	0.945(24)(4)(3)	0.939	0.050(9)(1)(3)	0.033(16)(4)(2)
Λ	1.103(13)(9)(3)	1.116	0.028(4)(1)(2)	0.144(15)(10)(2)
Σ	1.182(11)(2)(6)	1.193	0.0212(27)(1)(17)	0.187(15)(3)(4)
Ξ	1.301(12)(9)(1)	1.318	0.0100(10)(0)(4)	0.244(15)(12)(2)

$$\bar{\sigma}_{Bq} = (m_q/M_B)\partial M_B/\partial m_q$$

N. B. Masses are absolute calculations based upon heavy quark potential, which involves no chiral physics

Young & Thomas, arXiv:0901.3559 [nucl-th]



Sigma Commutator

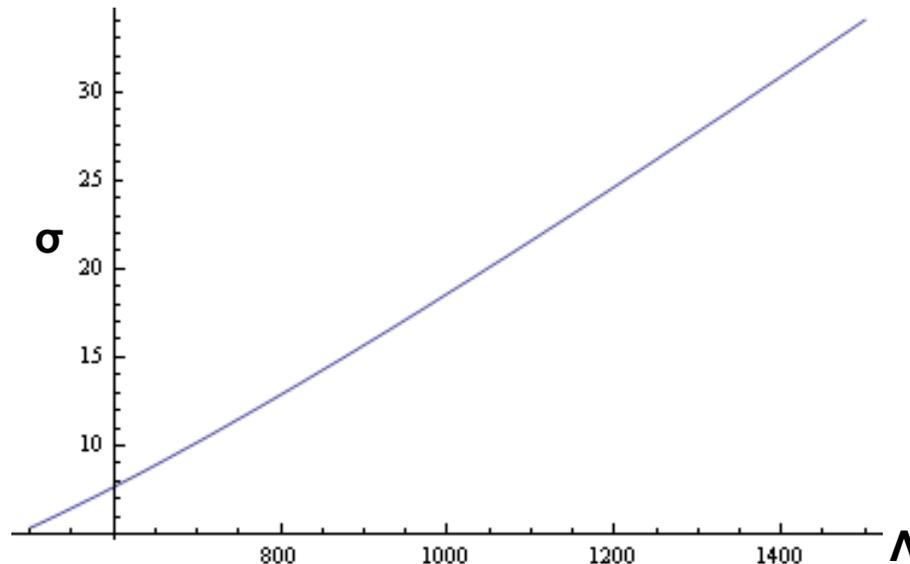
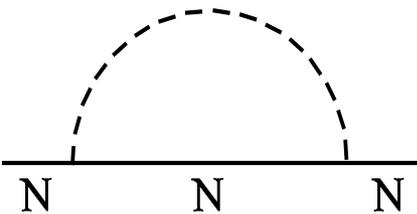
$$\sigma = \langle N | (m_u + m_d) (\bar{u} u + \bar{d} d)/2 | N \rangle \equiv m_q \partial M_N / \partial m_q$$

$$= \langle N | [Q_5, [Q_5, H_{\text{QCD}}]] | N \rangle$$

$$= \sigma_{\text{val}} + \sigma_{\text{sea}}$$

LNA

$$\delta\sigma = 35 \Lambda - 23 + \frac{9.6}{\Lambda} - \frac{3}{\Lambda^2} + \frac{0.8}{\Lambda^3} + \dots = 18 \text{ MeV} (\Lambda = 1 \text{ GeV})$$



Naïve Expansion Traditionally Used to Extract σ Terms is Hopeless!

- Leading-order expansion $O(1)$

$$M_N = M_0 + 2(\alpha_M + \beta_M)m_q + 2\sigma_M(2m_q + m_s)$$

$$M_\Lambda = M_0 + (\alpha_M + 2\beta_M)m_q + \alpha_M m_s + 2\sigma_M(2m_q + m_s)$$

$$M_\Sigma = M_0 + \frac{1}{3}(5\alpha_M + 2\beta_M)m_q + \frac{1}{3}(\alpha_M + 4\beta_M)m_s + 2\sigma_M(2m_q + m_s)$$

$$M_\Xi = M_0 + \frac{1}{3}(\alpha_M + 4\beta_M)m_q + \frac{1}{3}(5\alpha_M + 2\beta_M)m_s + 2\sigma_M(2m_q + m_s)$$

Need $O(m_\pi^6)$ to get accurate light quark σ term

While for strange condensate expansion is useless !

BUT through FRR have closed expression and can evaluate

Summary of Results of Combined Fits

(of 2008 LHPC & PACS-CS data)

B	Mass (GeV)	Expt.	$\bar{\sigma}_{Bl}$	$\bar{\sigma}_{Bs}$
N	0.945(24)(4)(3)	0.939	0.050(9)(1)(3)	0.033(16)(4)(2)
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$$\bar{\sigma}_{Bq} = (m_q/M_B) \partial M_B / \partial m_q$$

Of particular interest:

σ commutator well determined : $\sigma_{\pi N} = 47 (9) (1) (3) \text{ MeV}$

and strangeness sigma commutator small

$m_s \partial M_N / \partial m_s = 31 (15) (4) (2) \text{ MeV}$

NOT several 100 MeV !

Profound Consequences for Dark Matter Searches



Hadronic Uncertainties in the Elastic Scattering of Supersymmetric Dark Matter

John Ellis,^{1,*} Keith A. Olive,^{2,†} and Christopher Savage^{2,‡}

CERN-PH-TH/2008-005

UMN-TH-2631/08

FTPI-MINN-08/02

We find that the spin-independent cross section may vary by almost an order of magnitude for $48 \text{ MeV} < \Sigma_{\pi N} < 80 \text{ MeV}$, the $\pm 2\text{-}\sigma$ range according to the uncertainties in Table I. This uncertainty is already impacting the interpretations of experimental searches for cold dark matter. Propagating the $\pm 2\text{-}\sigma$ uncertainties in $\Delta_s^{(p)}$, the next most important parameter, we find a variation by a factor ~ 2 in the spin-dependent cross section. Since the spin-independent cross section may now be on the verge of detectability in certain models, and the uncertainty in the cross section is far greater, *we appeal for a greater, dedicated effort to reduce the experimental uncertainty in the π -nucleon σ term $\Sigma_{\pi N}$.* This quantity is not just an object of curiosity for those interested in the structure of the nucleon and non-perturbative strong-interaction effects: it may also be key to understanding new physics beyond the Standard Model.

$$\mathcal{L} = \alpha_{2i} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q}_i \gamma_\mu \gamma^5 q_i + \alpha_{3i} \bar{\chi} \chi \bar{q}_i q_i$$

spin
 σ terms

Neutralino (0.3 GeV / cc : WMAP)

Summary

- Strange content of N small
 - Less than 5% of μ^p and $\langle r^2 \rangle_{ch}^p$
- Theory agrees well but order of magnitude more accurate
- Major success of QCD : direct insight into “disconnected diagrams”
- analogue of Lamb shift

Hydrogen Atom, Electron (g-2)-factor, QED

$$g_e = 2 \left(1 + \frac{\alpha}{2\pi} - 0.328 \frac{\alpha^2}{\pi^2} + \dots \right)$$

Summary - 2

- **With lattice scale set by heavy quark (non-chiral) physics, agreement with octet masses is good at the 1-2% level!**
- **Strangeness content (condensate) is roughly an order of magnitude smaller than naively assumed**
- **Strangeness term usually dominates estimates of dark matter cross section - it should NOT!**

Summary - 3

FRR provides a very natural explanation of all of these results

i.e. The contributions of Goldstone boson (π and K and η) degrees of freedom is suppressed once the Compton wavelength is too small (<0.5 fm)

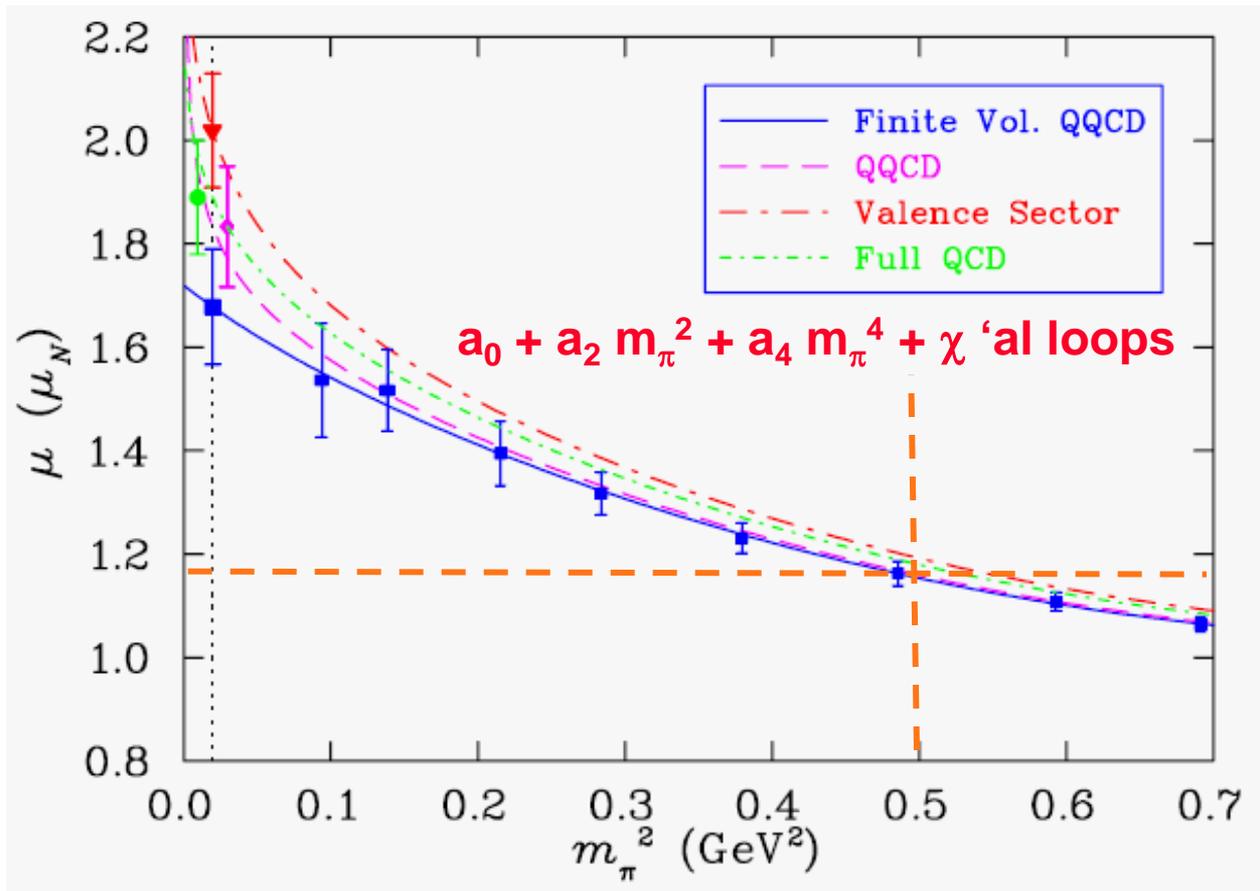
This corresponds to Goldstone boson masses of order 0.4 GeV and higher, which is where the physical K and η are located

Hence their contribution and the role of virtual strange quarks is suppressed





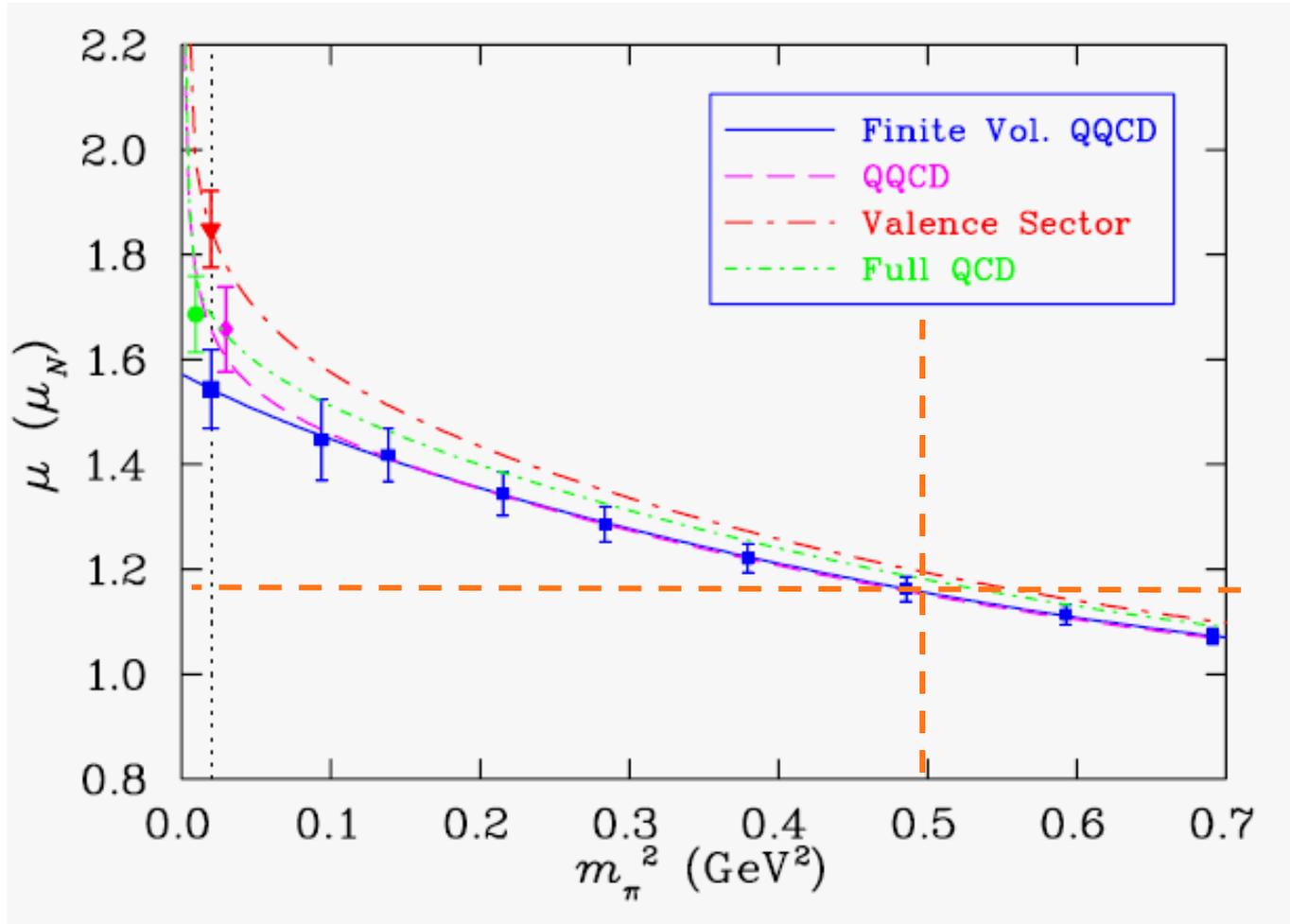
u^p_{valence} : QQCD Data Corrected for Full QCD Chiral Coefficients



← c.f. CQM
2/3 940/540
» 1.18

Lattice data from Zanotti et al. ; Chiral analysis Leinweber et al.

u^Σ valence



← Universal Here!

Convergence LNA to NLNA Again Excellent (Effect of Decuplet)

