

# Importance of $e+D$ scattering at a collider

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“4<sup>th</sup> workshop on the LHeC”

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# Overview

- Why the deuteron?
- Nuclear effects from low to high  $x$ 
  - beyond obvious LHeC strength
- Proton/neutron tagging
- Diffraction
  - Only touched upon in this talk

# Why deuteron?

## Deuteron as effective neutron beam

## Quark flavor decomposition

$$F_2(p) \propto 4u + d$$

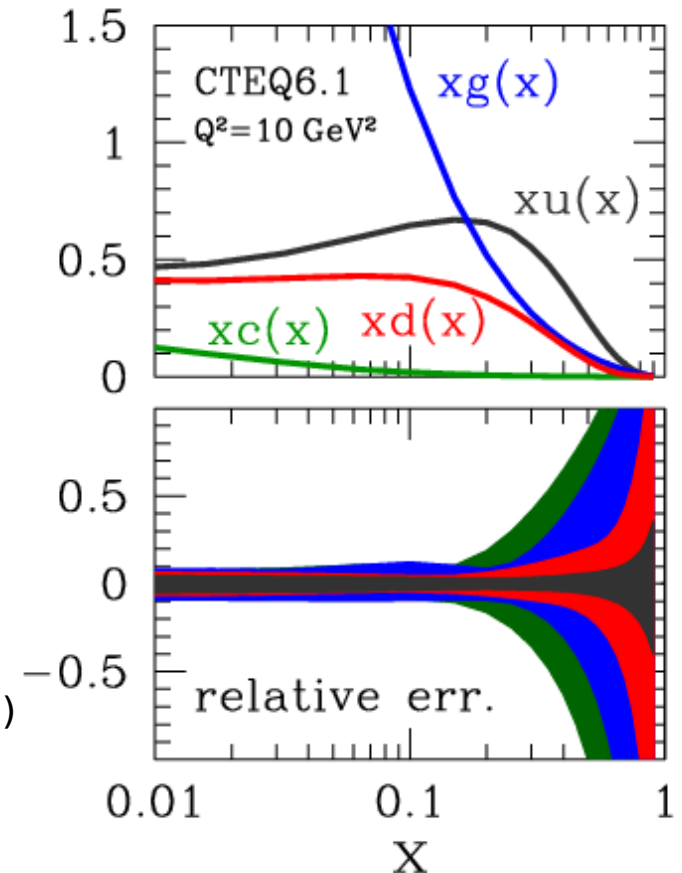
$$F_2(n) \propto u + 4d$$

## Particularly important at large $x$

- Large d-quark uncertainty
- $d/u$  ratio at  $x \rightarrow 1$  probes non perturbative proton structure

Accardi et al. [CTEQ-JLab collab.] PRD84(2011)

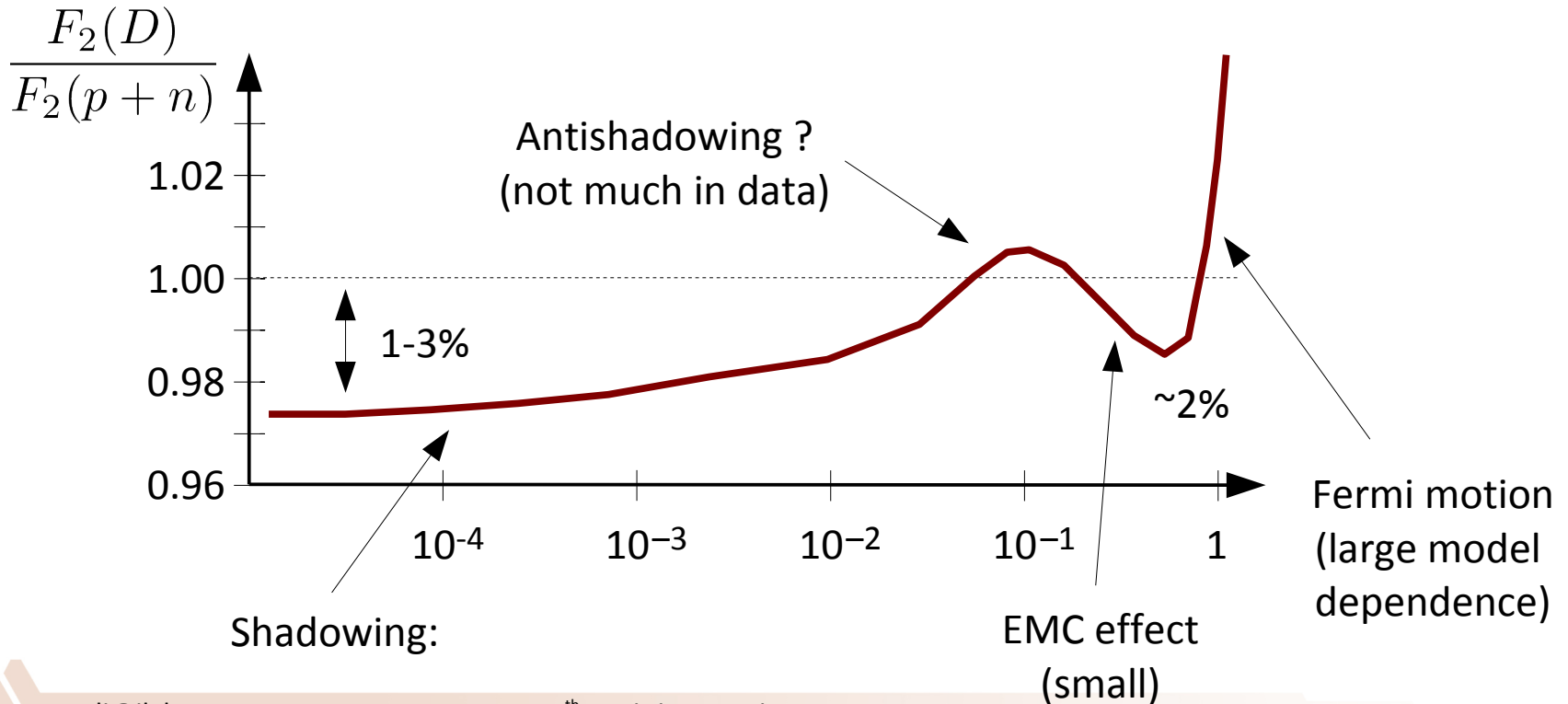
## At $x \lesssim 10^{-2}$ sea quarks dominate, expect $F_2(p) \approx F_2(n)$



# Why deuteron?

## □ As baseline for nuclear PDFs

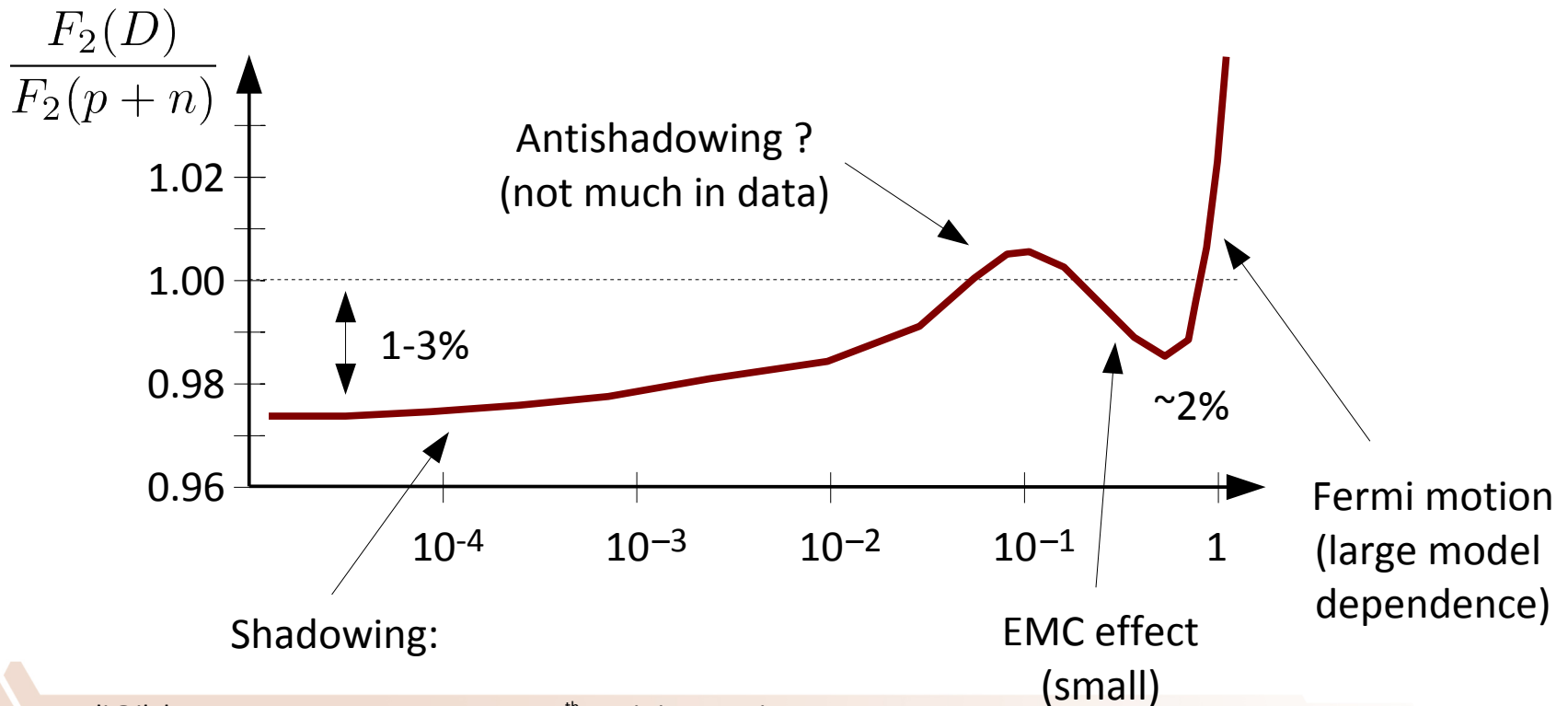
- $F_2(A)/F_2(D)$  is measured: approx. correct for isospin, remove syst.
  - Ideally,  $F_2(A)/[ZF_2(p) + (A - Z)F_2(n)]$
- But  $F_2(D) \neq F_2(p) + F_2(n)$  , need to measure it



# Why deuteron?

## □ The simplest nucleus:

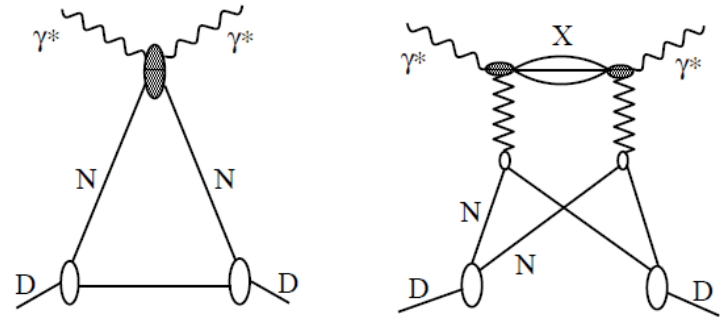
- Nuclear few-body calculations of  $p, n$  wave function available
- Testbed for nuclear effects calculations and modeling



# Nuclear effects 1 - shadowing

## □ Nuclear shadowing

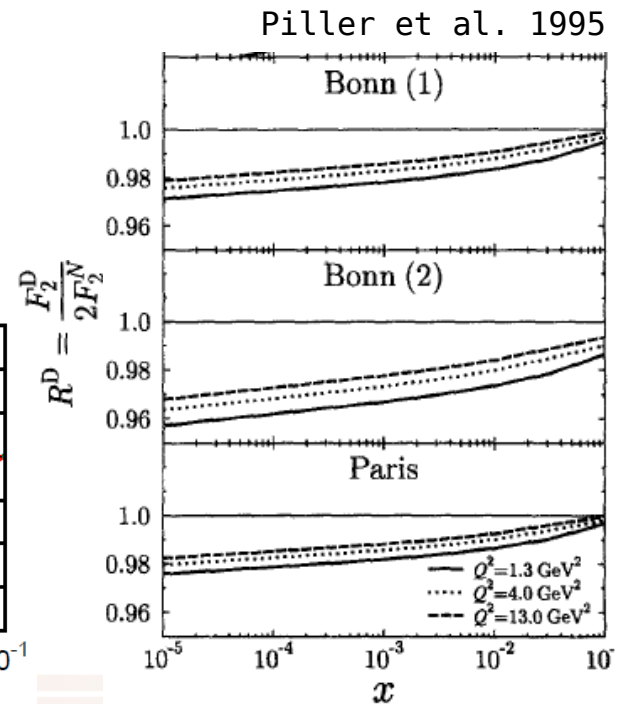
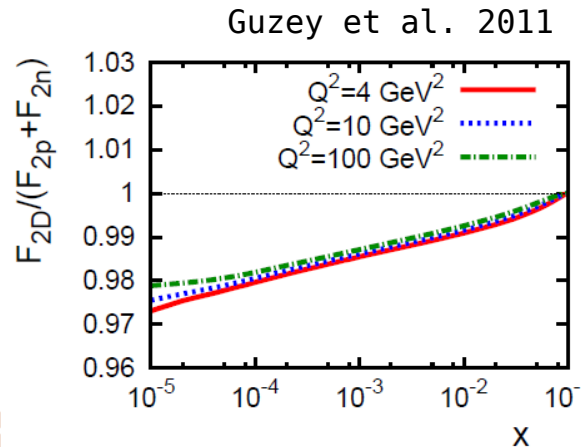
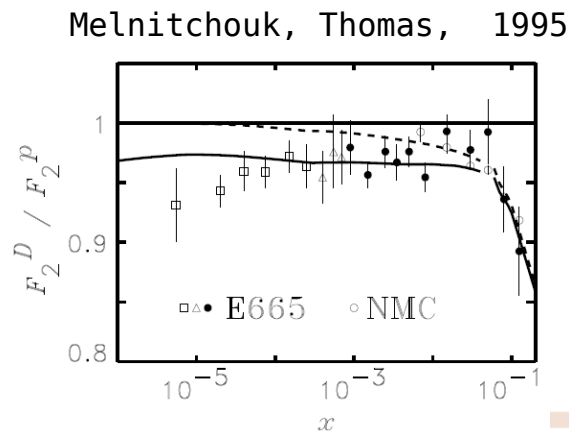
- Double scattering only
- At low  $x \ll 0.01$ , Glauber-Gribov, no model dependency
- Connection to diffraction (on  $p$  and  $n$ )



Badelek, Kwiecinsky, NPB370(1992)

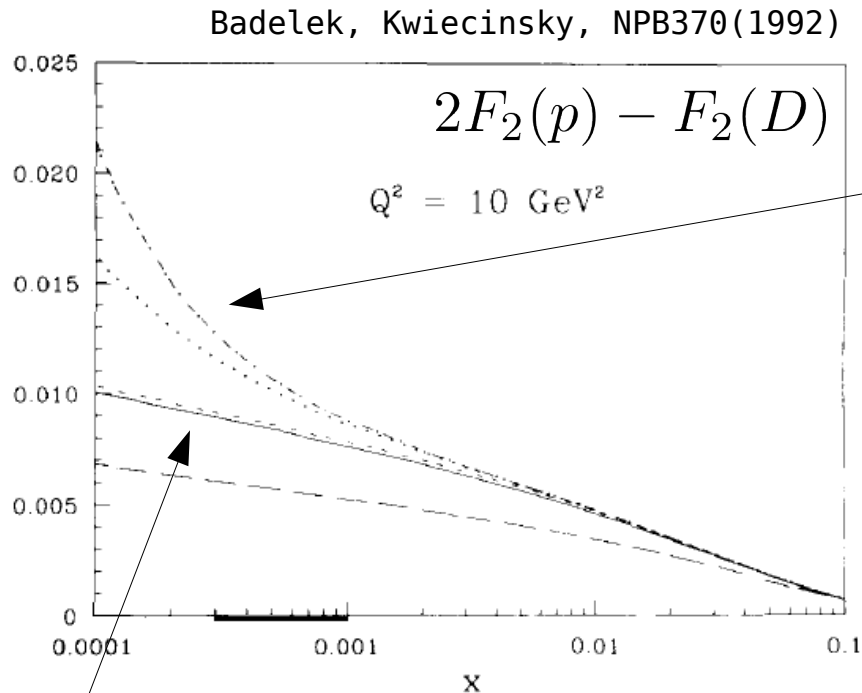
- All implementations: 1-3% shadowing at  $10^{-5} < x < 10^{-2}$

- Main uncertainty: deuteron wave function



# Nuclear effects 1 - shadowing

## Shadowing + gluon recombinations



With recombination  
(2 different gluon PDF)  
→ more shadowing

Without recombination

# Nuclear effects 1 - shadowing

## □ Shadowing + gluon recombinations

### Test shadowing calculations in controlled setting

→ approach to saturation

→ Access to diffractive  $F_2^D(n)$

Without recombination

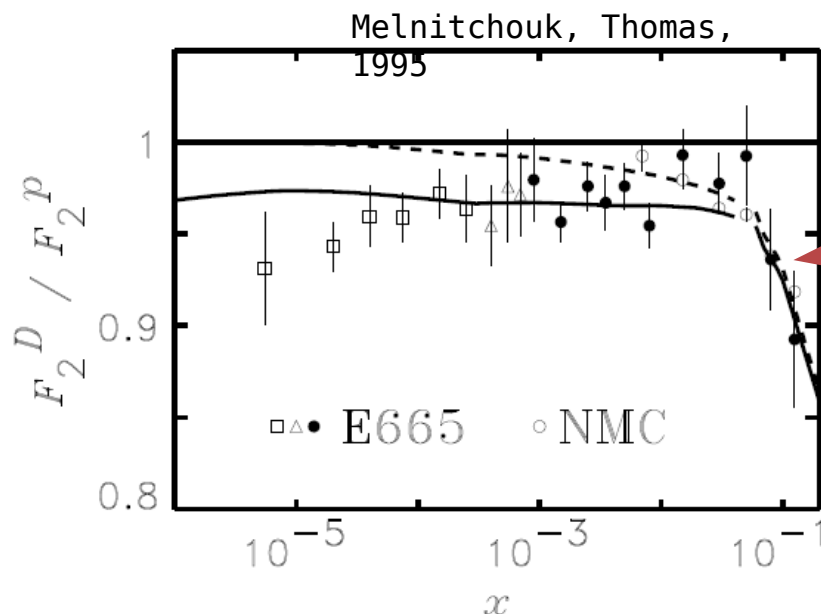
recombination  
(gluons)



# Nuclear effects 2 - antishadowing

## □ Antishadowing

- Not obvious in data because  $F_2(p) \neq F_2(n)$  at  $x \gtrsim 0.01$
- The baseline is not 1



## □ Needs direct measurements of $F_2(n)$

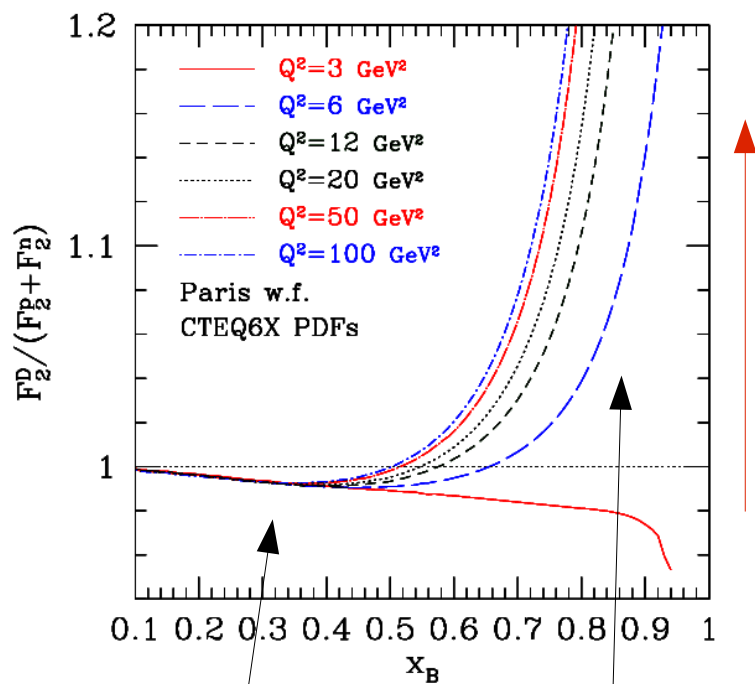
- Proton tagging, see later
- Data exist from BoNuS at  $x > 0.3$ , low  $Q$ ; lower  $x$ , higher  $Q$  at JLab12

# Nuclear effects 3 - EMC, Binding, Fermi motion

□ Origin of EMC effect still a mystery after 30 years:

- $x > 0.1$  is a complex region
- many theoretical uncertainties

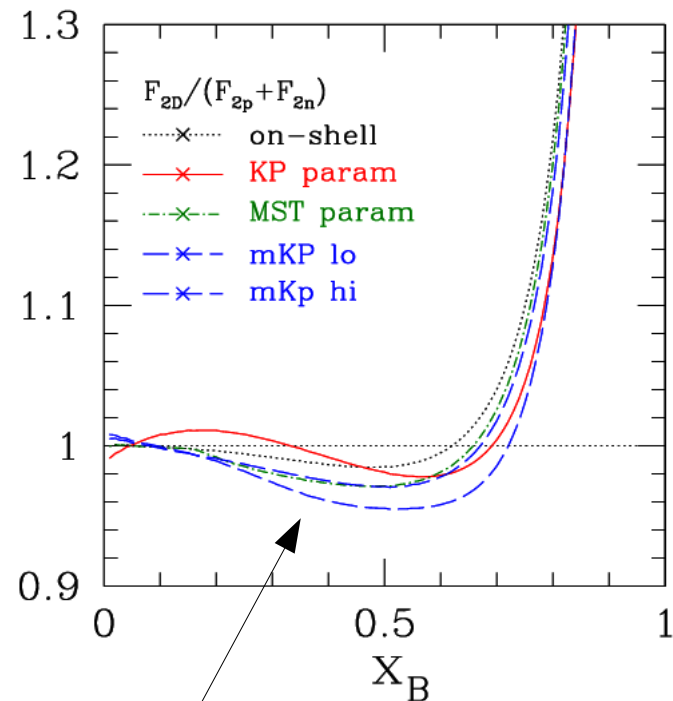
Accardi et al. [CTEQ-JLab collab.] PRD84(2011)



Binding

Fermi motion

1/Q<sup>2</sup> effects:  
 - Higher twists  
 - Target Mass  
 - p<sup>2</sup>/Q<sup>2</sup> effects

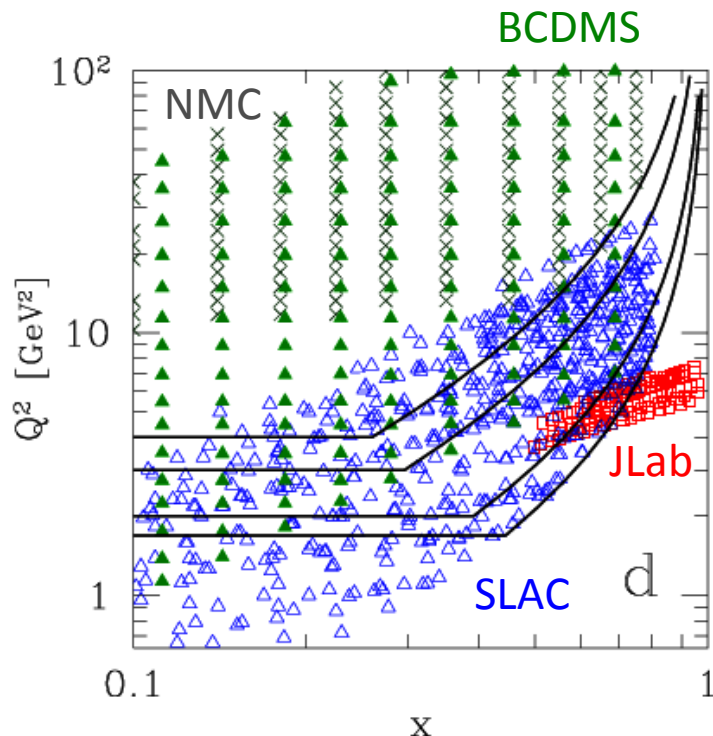


Nucleon off-shellness effects

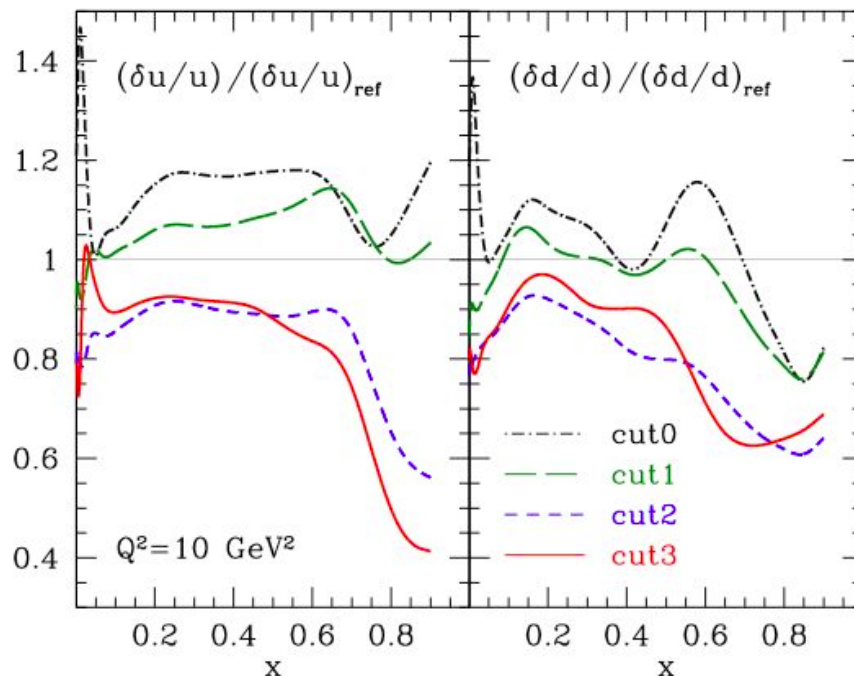
# Nuclear effects 3 - EMC, Binding, Fermi motion

## □ Important for:

- Fit of  $d$  quark at large  $x$  Accardi et al. [CTEQ-JLab] PRD84(2011)
- Constraining nucl. corrections, when compared to free proton data (for example,  $p+p \rightarrow W(Z)+X$ , PVDIS, ...) Accardi, ECT\*, May 2012

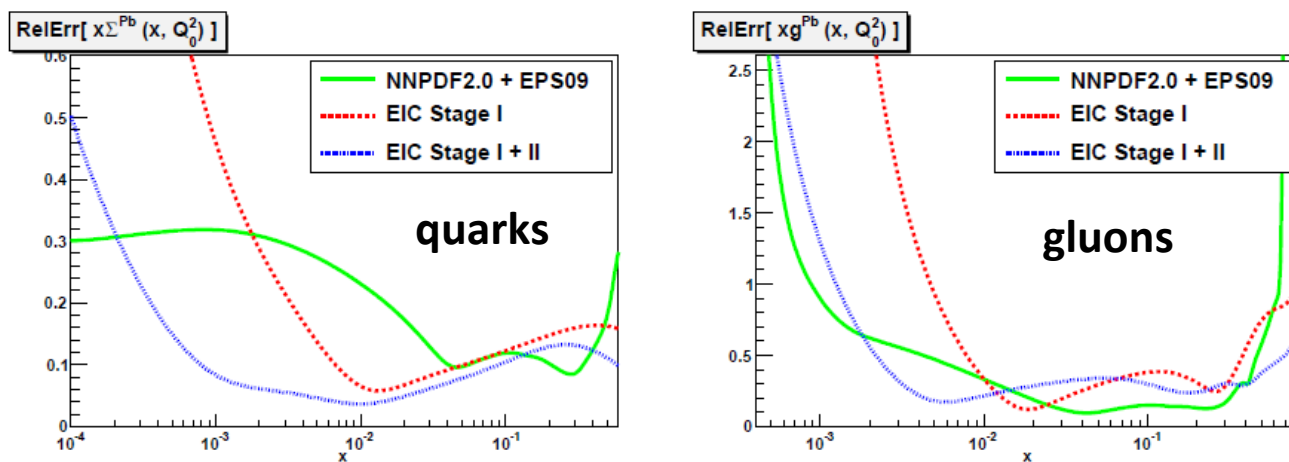


Accardi et al. [CTEQ-JLab] PRD81(2010)



# Nuclear effects 3 - EMC, Binding, Fermi motion

- Not impossible to have good data at a collider
  - For example at EIC, 1 year of data at  $L = 4 \text{ fb}^{-1}$  for 5+3 energies



Accardi, Guzey, Rojo, arXiv:1106.383

- Advantage: large  $Q^2$  leverage
  - Suppress  $1/Q^2$  power corrections
  - Gluon EMC effect via scaling violation of  $F_2$ 
    - This may be as revolutionary as original quark EMC effect

# Nuclear effects 3 - EMC, Binding, Fermi motion

- Not impossible to have good data at a collider

- Fermi motion

***d* quark at large  $x$**

**Test of nuclear corrections**

**Gluon EMC effect**

- Advantages

- Suppress  $1/Q^2$  power corrections

- Gluon EMC effect via scaling violation of  $F_2$

- This may be as revolutionary as original quark EMC effect

# Nuclear effects 4 - superfast quarks

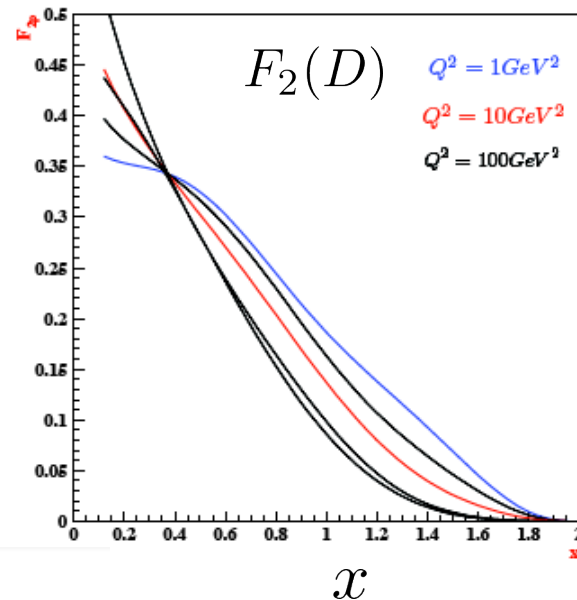
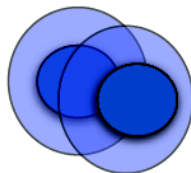
Sargsian et al JPG29 (2003)

□  $F_2(D)$  can go to  $x > 1$  at large  $Q^2$ :

– “superfast quarks”

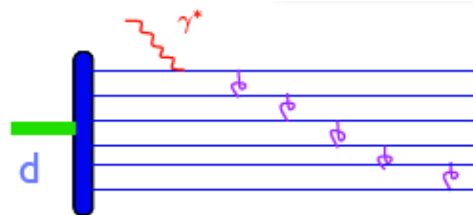
□ Scatter on faster-than-average nucleon

– Probe short-range  
 $NN$  potential



□ Scattering on exotica

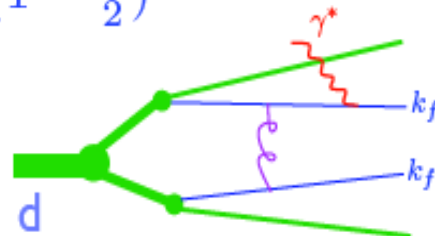
– e.g., 6-quark bags



$$F_{2,(6q)} \sim \left(1 - \frac{x}{2}\right)^{10}$$

□ Novel QCD mechanisms

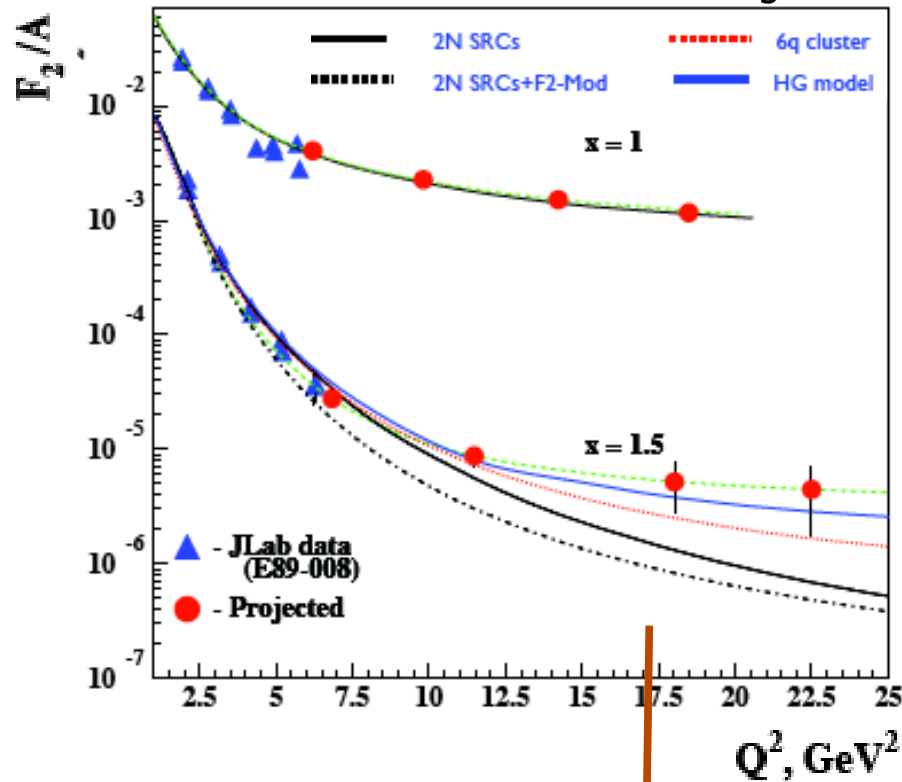
– Hard-gluon exchange?



# Nuclear effects 4 - superfast quarks

## Signatures

M. Sargsian

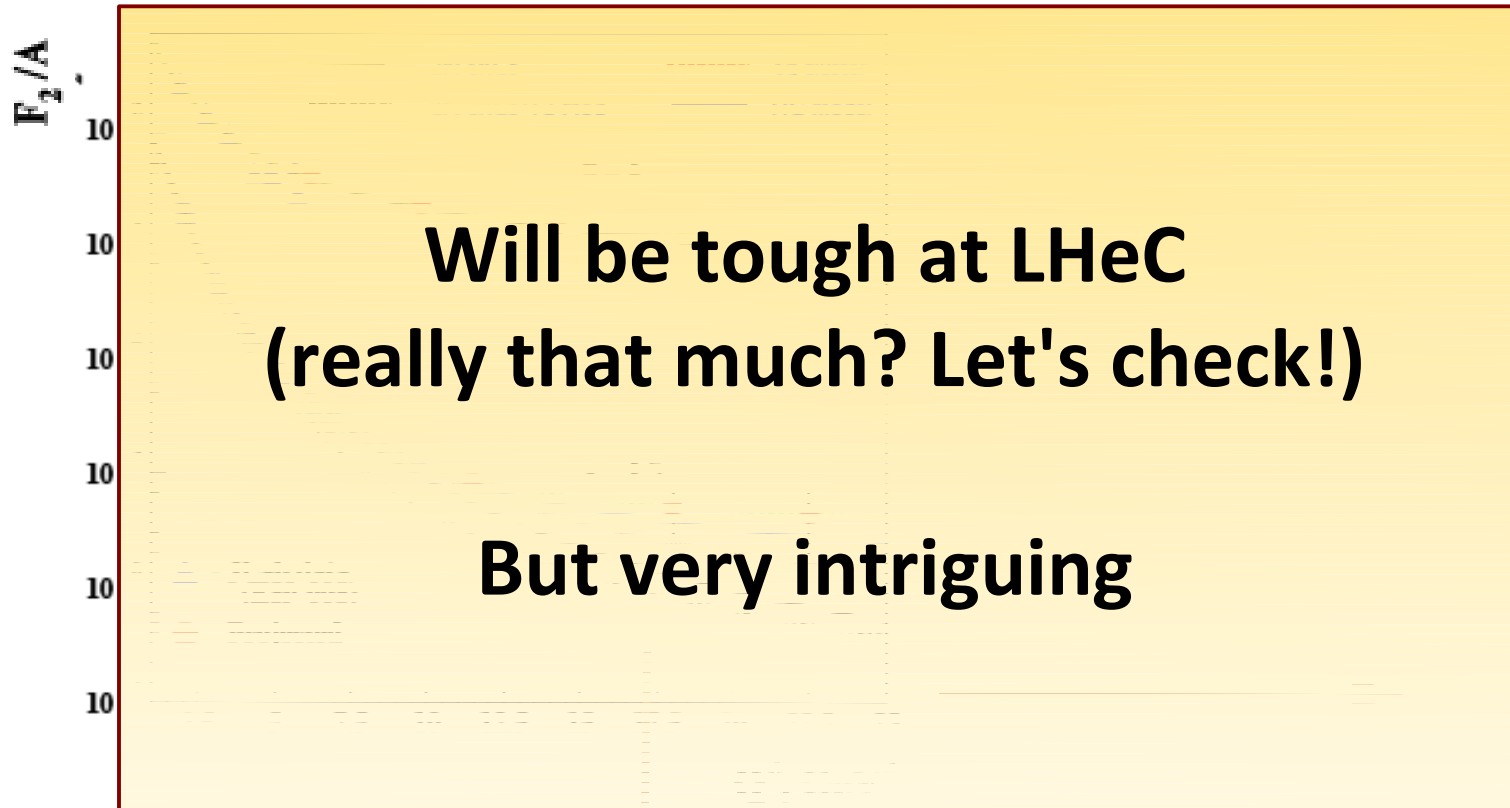


EIC, LHeC:

larger  $x$ , larger  $Q^2$

# Nuclear effects 4 - superfast quarks

## □ Signatures



**EIC, LHeC:  
larger x, larger  $Q^2$**

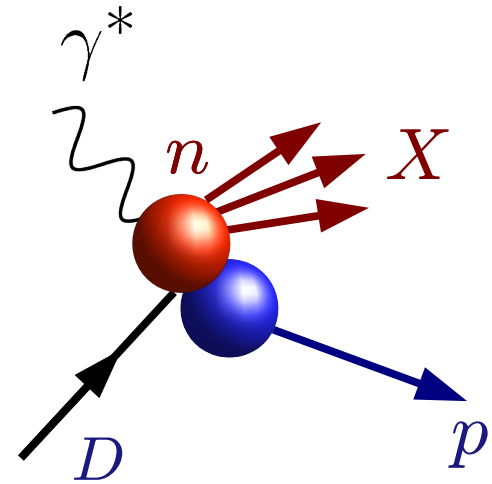


# Nucleon tagging

## Proton tagging: scattering on neutron

$$F_2 = F_2(x, Q^2, \vec{p})$$

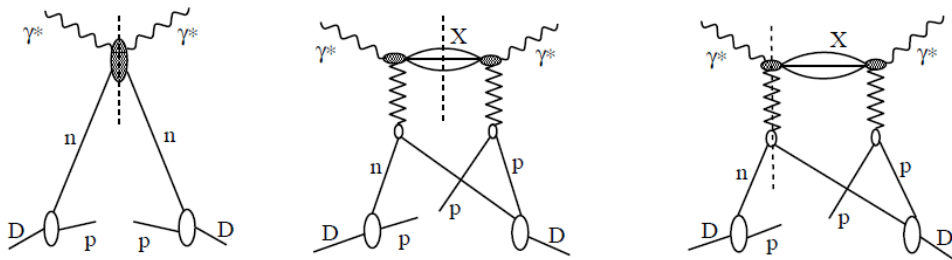
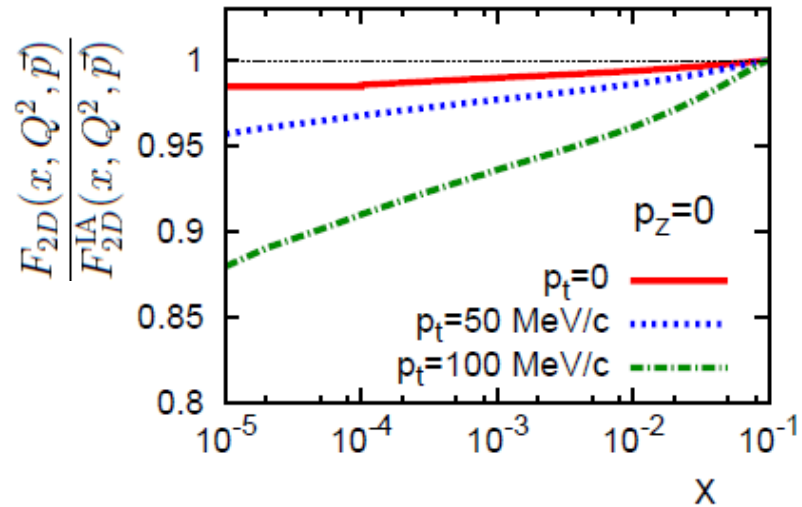
$\vec{p}$  = spectator momentum



– But: shadowing, binding, off-shellness, F. motion...

## Small x: shadowing corrections

- Minimized when  $\vec{p} \approx 0$
- Diffraction on “quasi-free neutrons”



# Nucleon tagging

## Small x: Final State Interactions

- Minimized for:
  - Spectators anti-parallel to  $\gamma^*$
  - Slow protons
- “quasi-free neutrons”

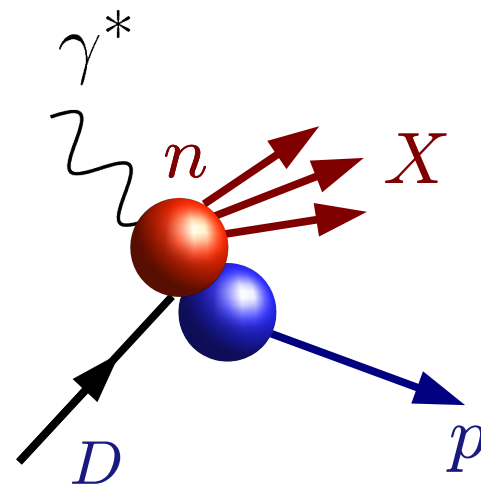
## Minimizing nuclear w.f. uncertainty

- Take suitable ratios of  $F_2$

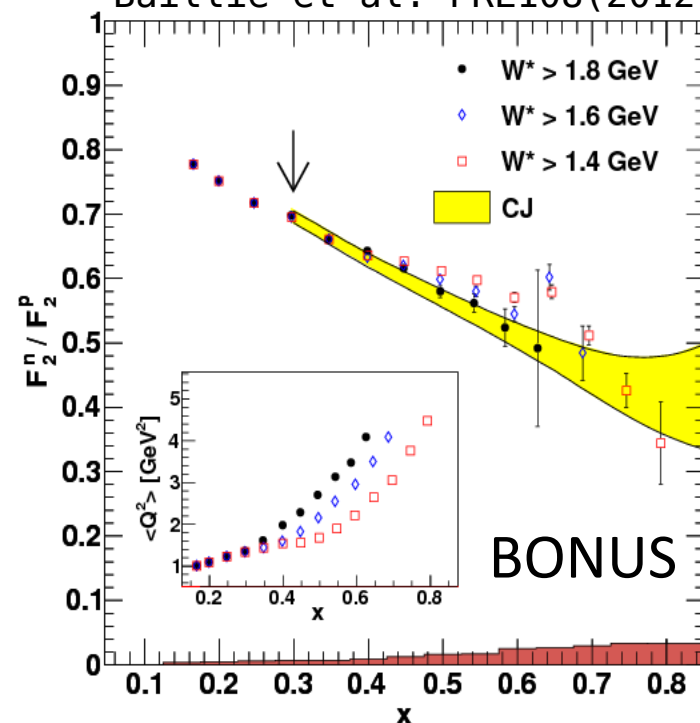
## Neutron tagging:

in-medium modifications

- By comparing tagged to free protons



Baillie et al. PRL108(2012)



# Nucleon

## Small

– M

– “

## Mini

– T

## Neut

in-m

– B

Collider ideal for nucleon tagging,  
especially neutron

**CHALLENGE:** good control and  
resolution of spectator momentum

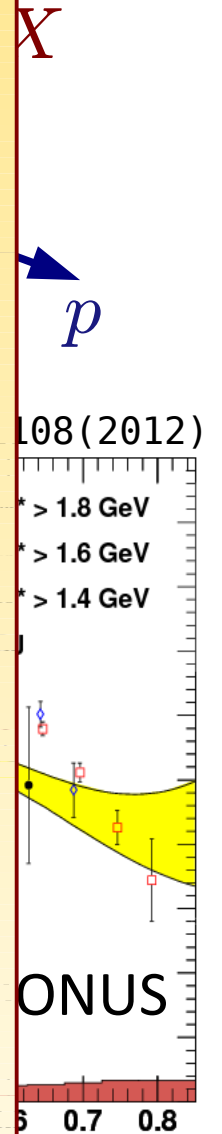
**BUT:**

Neutron DIS, diffraction

Baseline for antishad. in  $D$

Free vs. bound, off-shell protons

...



# Summary: why deuteron?

- Flavor separation, baseline for nuclear PDF fits
- Nuclear effects from low to high  $x$ 
  - Verify shadowing calculations, approach to sat.,  $F_2^D(n)$
  - Bound nucleons without  $1/Q^2$  corrections
  - “Superfast quarks”
- Proton/neutron tagging
  - DIS, diffraction on neutrons
  - Free vs. bound, off-shell protons
- Diffraction (not covered in this talk)
  - Coherent, breakup, incoherent